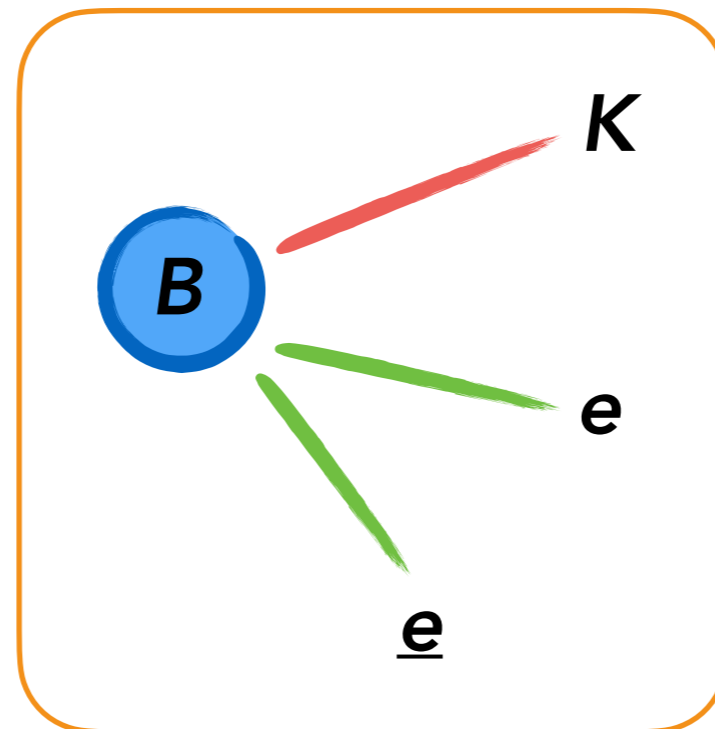


# EFFECTIVE FIELD THEORIES in PARTICLE PHYSICS

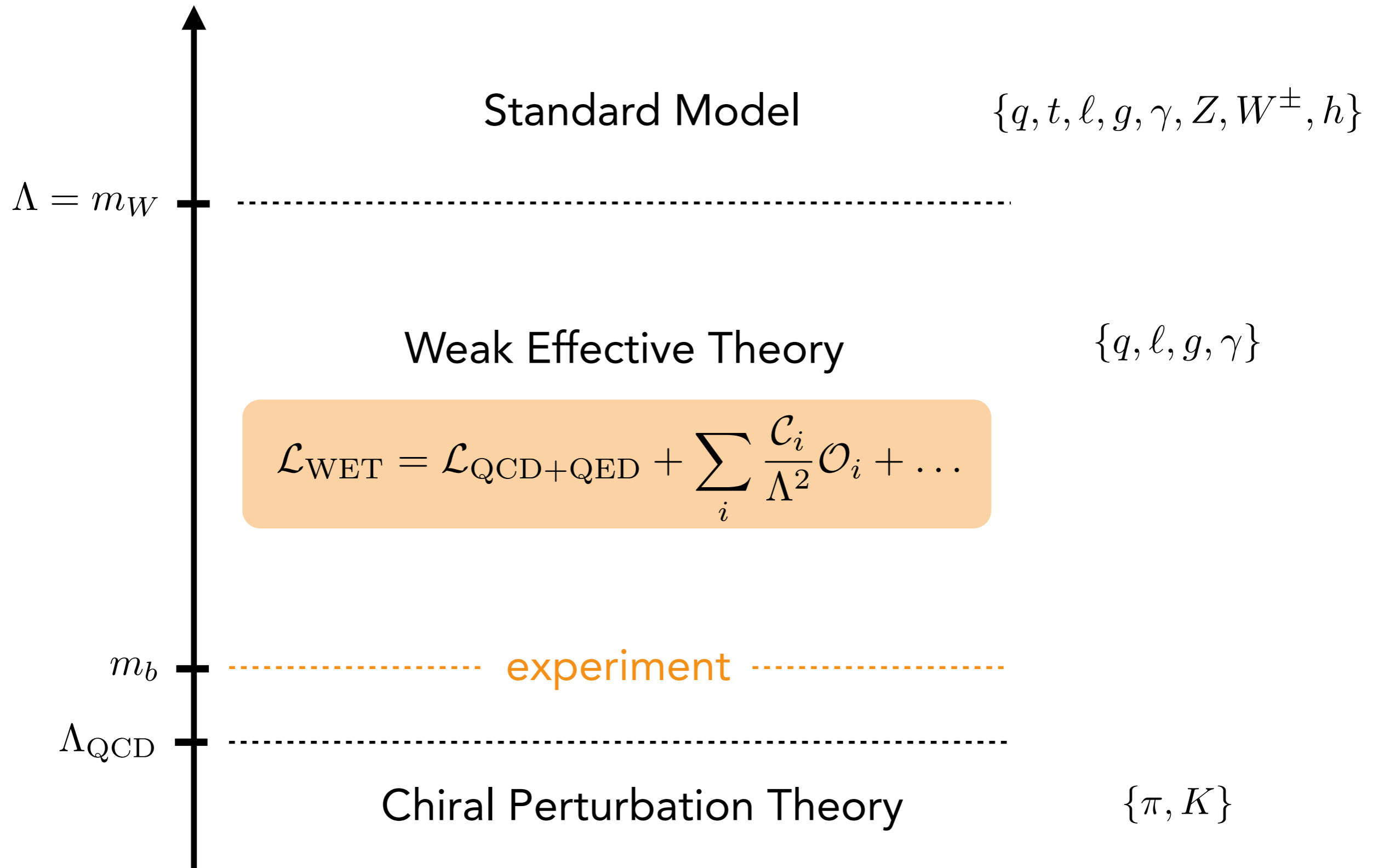
Susanne Westhoff  
Heidelberg University

# Part II

## Weak Effective Theory



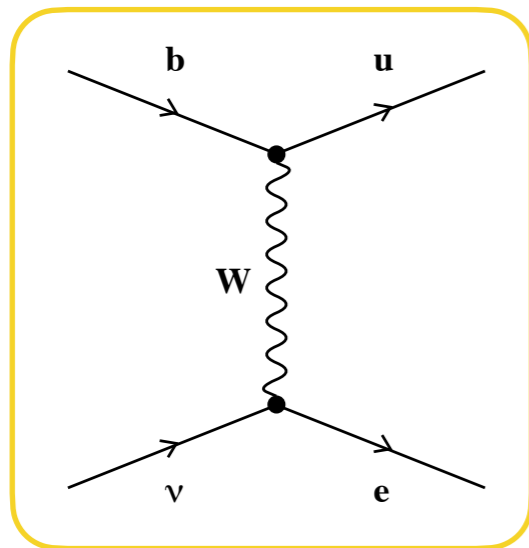
# Multi-scale physics



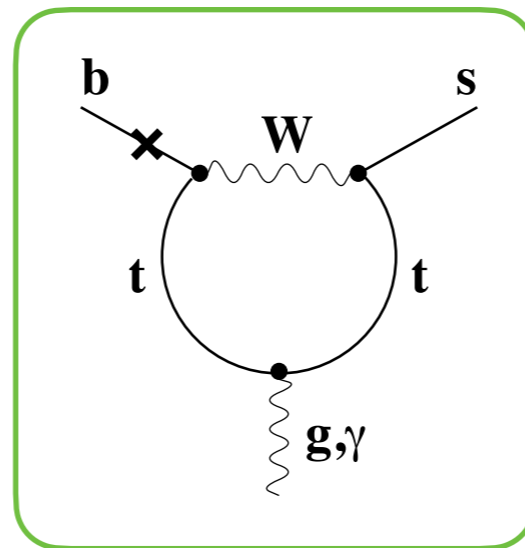
# Weak Effective Theory

Goal: precise description of flavor-changing processes.

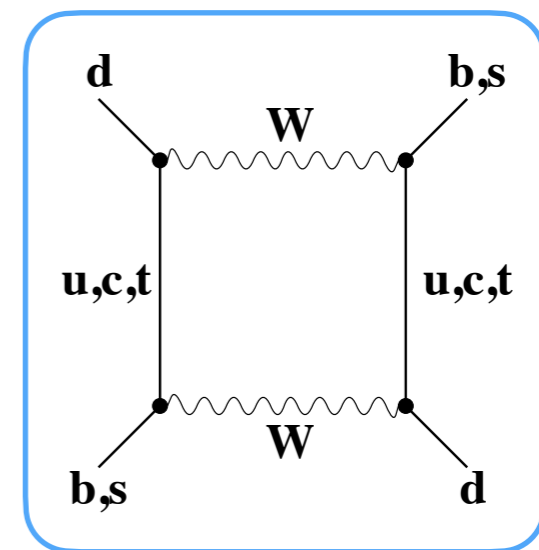
$$\bar{B}^0 \rightarrow \pi^+ e^- \bar{\nu}_e$$



$$B \rightarrow X_s \gamma$$



$$B \leftrightarrow \bar{B}$$



$$\mathcal{L}_{\text{WET}} = -\frac{4G_F}{\sqrt{2}} \sum_i C_i O_i^{(6)} + \mathcal{O}\left(\frac{m_b^2}{m_W^2}\right)$$

What are the relevant operators?

# Matching

$$\mathcal{M} = \langle f | \mathcal{L}_{\text{SM}} | i \rangle \stackrel{!}{=} -\frac{4G_F}{\sqrt{2}} \sum_i C_i(\mu) \langle f | O_i^{(6)}(\mu) | i \rangle$$

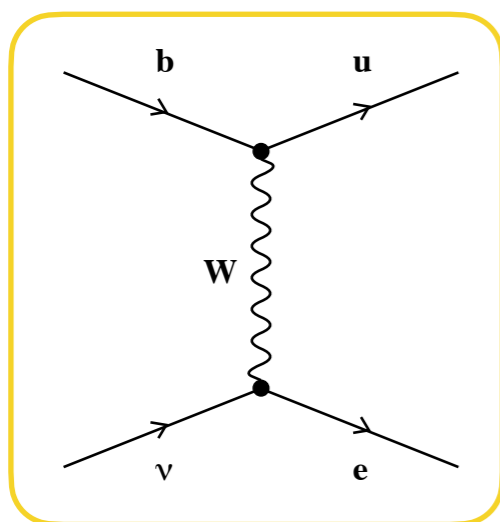
- match full SM amplitude onto weak effective theory
- include all operators that respect the SM symmetries

# Matching

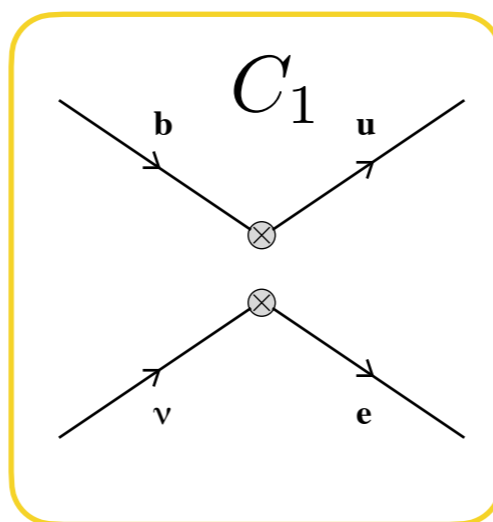
$$\mathcal{M} = \langle f | \mathcal{L}_{\text{SM}} | i \rangle \stackrel{!}{=} -\frac{4G_F}{\sqrt{2}} \sum_i C_i(\mu) \langle f | O_i^{(6)}(\mu) | i \rangle$$

- match full SM amplitude onto weak effective theory
- include all operators that respect the SM symmetries

$$\overline{B}^0 \rightarrow \pi^+ e^- \bar{\nu}_e$$



match  
 $\longleftrightarrow$



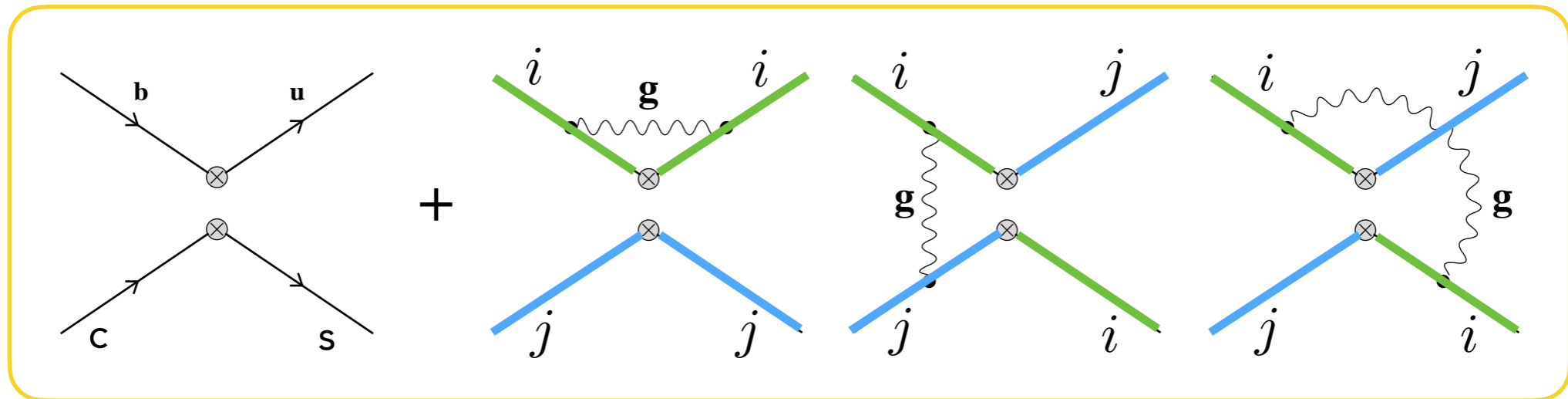
$$\mathcal{L}_{\text{eff}} = -\frac{4G_F}{\sqrt{2}} C_1 O_1$$

$$O_1 = V_{ub} (\bar{u}_L \gamma^\mu b_L) (\bar{e}_L \gamma_\mu \nu_L)$$

$$C_1 = 1$$

# Matching with QCD effects

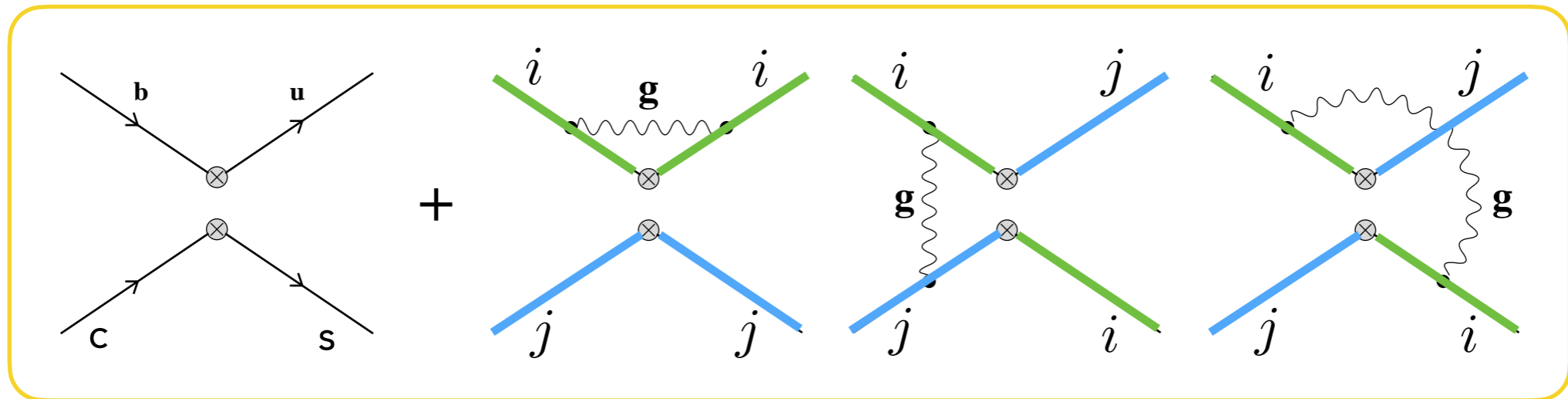
$$\overline{B}^0 \rightarrow \pi^+ D_s^-$$



$$\mathcal{L}_{\text{eff}} = -\frac{4G_F}{\sqrt{2}} V_{cs}^* V_{ub} \left[ C_1(\mu) \bar{s}_L^j \gamma_\mu c_L^j \bar{u}_L^i \gamma^\mu b_L^i + C_2(\mu) \bar{s}_L^i \gamma_\mu c_L^j \bar{u}_L^j \gamma^\mu b_L^i \right]$$

# Matching with QCD effects

$$\overline{B}^0 \rightarrow \pi^+ D_s^-$$



$$\mathcal{L}_{\text{eff}} = -\frac{4G_F}{\sqrt{2}} V_{cs}^* V_{ub} \left[ C_1(\mu) \bar{s}_L^j \gamma_\mu c_L^j \bar{u}_L^i \gamma^\mu b_L^i + C_2(\mu) \bar{s}_L^i \gamma_\mu c_L^j \bar{u}_L^j \gamma^\mu b_L^i \right]$$

- Wilson coefficients depend on matching scale:

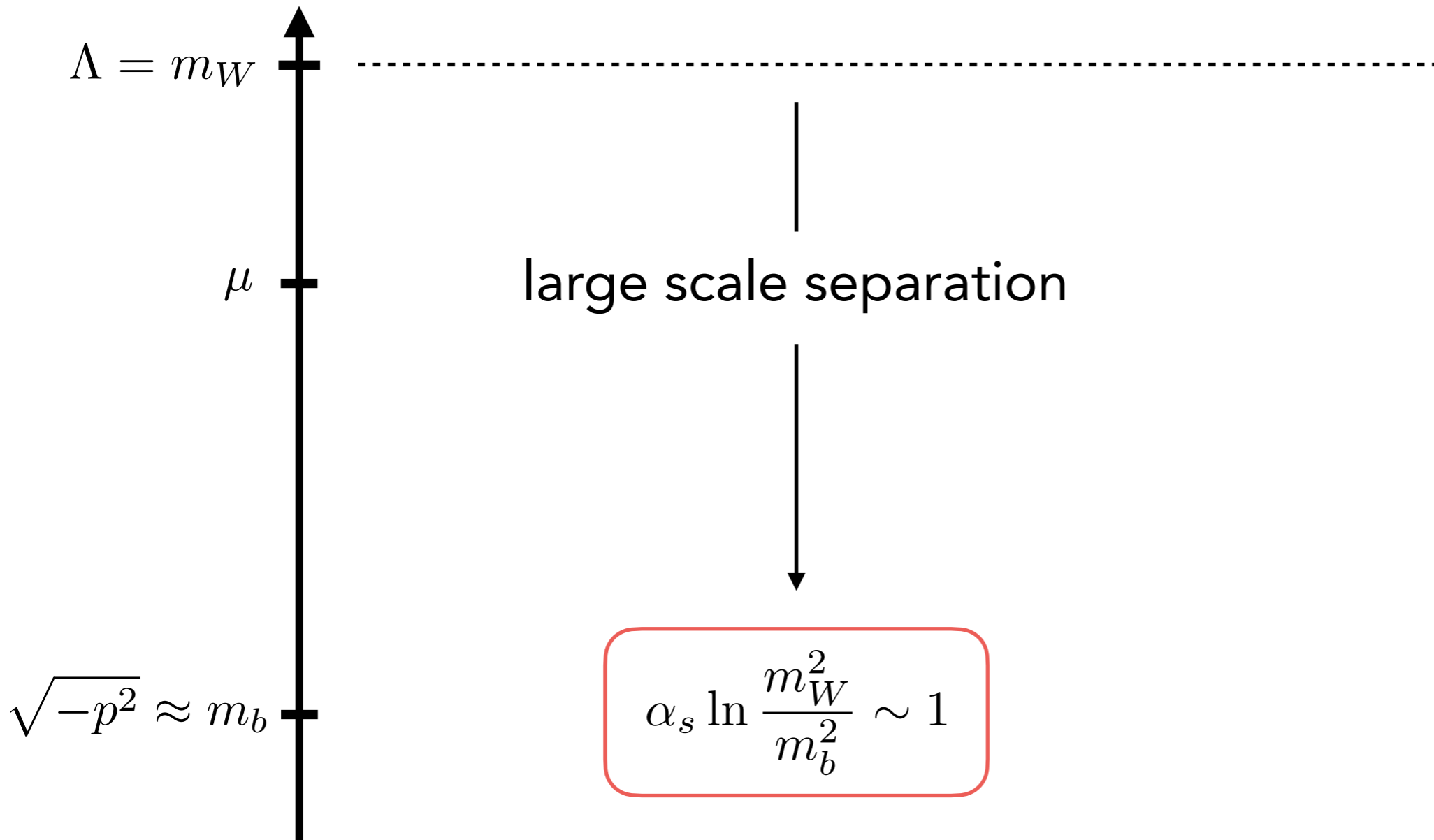
$$C_1(\mu) = 1 + \frac{3}{N_c} \frac{\alpha_s(\mu)}{4\pi} \left( \ln \frac{M_W^2}{\mu^2} - \frac{11}{6} \right) + O(\alpha_s^2)$$

$$C_2(\mu) = -3 \frac{\alpha_s(\mu)}{4\pi} \left( \ln \frac{M_W^2}{\mu^2} - \frac{11}{6} \right) + O(\alpha_s^2)$$



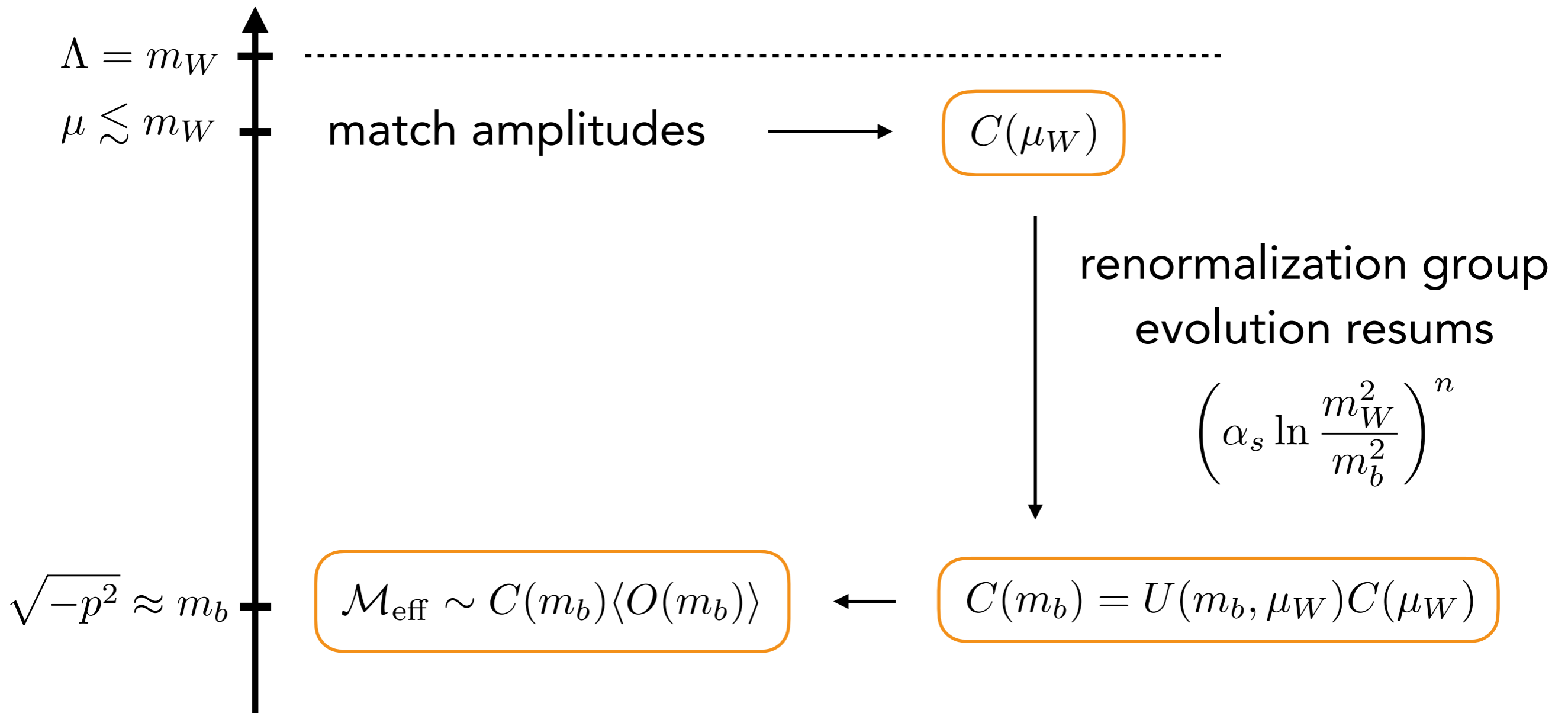
# Running couplings

factorize amplitude:  $1 + \alpha_s \ln \frac{M_W^2}{-p^2} = \underbrace{\left(1 + \alpha_s \ln \frac{M_W^2}{\mu^2}\right)}_{C(\mu)} \underbrace{\left(1 + \alpha_s \ln \frac{\mu^2}{-p^2}\right)}_{\langle O(\mu) \rangle} + \dots$



# Running couplings

factorize amplitude:  $1 + \alpha_s \ln \frac{M_W^2}{-p^2} = \underbrace{\left(1 + \alpha_s \ln \frac{M_W^2}{\mu^2}\right)}_{C(\mu)} \underbrace{\left(1 + \alpha_s \ln \frac{\mu^2}{-p^2}\right)}_{\langle O(\mu) \rangle} + \dots$



# Running couplings - example

$$\bar{B}^0 \rightarrow \pi^+ D_s^-$$

$$\mathcal{L}_{\text{eff}} = -\frac{4G_F}{\sqrt{2}} V_{cs}^* V_{ub} [C_1(\mu) \bar{s}_L^j \gamma_\mu c_L^j \bar{u}_L^i \gamma^\mu b_L^i + C_2(\mu) \bar{s}_L^i \gamma_\mu c_L^j \bar{u}_L^j \gamma^\mu b_L^i]$$

$$C_1(m_W) = 1 + \mathcal{O}(\alpha_s)$$

$$C_2(m_W) = 0 + \mathcal{O}(\alpha_s)$$

match & run



$$C_1(m_b) \approx 1.1$$

$$C_2(m_b) \approx -0.2$$

Effect on decay branching ratio  $\mathcal{B}(\bar{B}^0 \rightarrow \pi^+ D_s^-) \sim C_i^2$ .

# Summary Part II

## Weak Effective Theory:

- good description of charged currents at  $E \ll m_W$
- effective Lagrangian:

$$\mathcal{L}_{\text{WET}} = -\frac{4G_F}{\sqrt{2}} \sum_i C_i O_i^{(6)} + \mathcal{O}\left(\frac{m_b^2}{m_W^2}\right)$$

Obtain  $B$ -physics amplitude  $\mathcal{M}_{\text{eff}} \sim C(m_b) \langle O(m_b) \rangle$  :

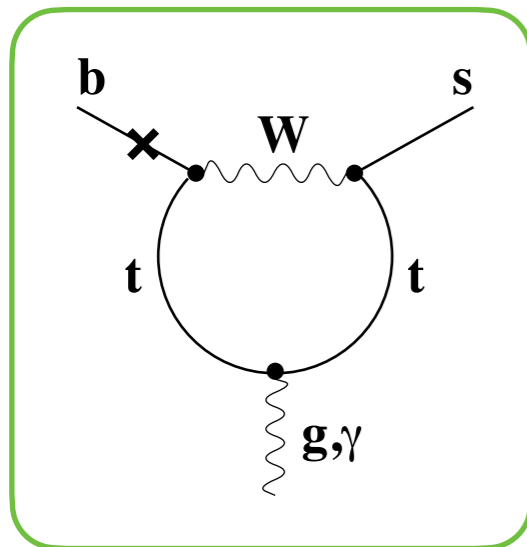
- match  $\longrightarrow C(\mu_W)$  (flavor-universal)
- evolve  $\longrightarrow C(m_b) = U(m_b, \mu_W) C(\mu_W)$

WET makes more precise predictions than full theory.\*

\*at the same perturbative order

# Your turn

$$B \rightarrow X_s \gamma$$

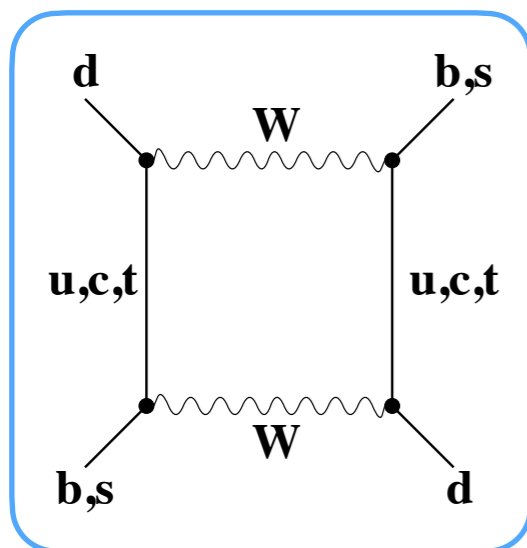


$$Q_{7\gamma} = \frac{e}{8\pi^2} m_b \bar{s}_i \sigma^{\mu\nu} (1 + \gamma_5) b_i F_{\mu\nu}$$

$$Q_{8G} = \frac{g}{8\pi^2} m_b \bar{s}_i \sigma^{\mu\nu} (1 + \gamma_5) T_{ij}^a b_j G_{\mu\nu}^a$$

- Draw Feynman diagrams of the two operators.
- Why does  $Q_{8G}$  contribute to  $B \rightarrow X_s \gamma$ ?

$$B \leftrightarrow \bar{B}$$



- Which WET operators contribute?